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One-dimensional percolation problem with further neighbour bonds—transfer-matrix approach

Z Q Zhang, F C Pu and T C Li

Institute of Physics, Chinese Academy of Sciences, Beijing, China

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Abstract. The transfer-matrix method is used to find the exact expressions for the correlation length in the critical region for both one-dimensional site and bond percolation problems with bonds connecting L th-nearest neighbours (for any finite L). For the site percolation, the correlation length exponent ν is found to be L , consistent with the result obtained from the generating function method. For the bond percolation, we find $\nu = L(L + 1)/2$.

1. Introduction

In the percolation problem, the one-dimensional system is one of the few cases where exact solutions can be found (for reviews see Stauffer 1979 and Essam 1980). One-dimensional site percolation with bonds connecting L th-nearest neighbour sites has been solved by the generating function method (Klein *et al* 1978). Although the critical occupation probability p_c is trivial in one dimension (p_c is always equal to 1), the critical exponents are found to be L dependent. In particular, the correlation length exponent ν is found to be equal to L . Such a 'bond range' dependence of the critical behaviour is closely related to the corresponding 'thermal' problem with multi-spin interactions (Klein *et al* 1978).

Recently, the transfer-matrix method has been used to find the exact critical behaviour for both one-dimensional site and bond percolation systems with further neighbour bonds (Zhang and Shen 1982). In that work, the authors showed that the transfer-matrix method gives the consistent result $\nu = L$ in the site percolation case. For the bond percolation, much richer critical behaviour was found. If all the L th-nearest neighbour bonds have equal occupation probability, ν is found to be $L(L + 1)/2$. However, in that work, the calculations are only done for small L (L up to 3); no proof has been given for the case of general L . In this work, we give a systematic extension of the transfer-matrix method to one-dimensional percolation systems from previous low L values to any finite L .

2. Transfer matrix

Consider a linear chain with bonds connecting L th-nearest neighbours (cf figure 1). Following the method used by Zhang and Shen (1982), we divide the chain into overlapping columns each containing L sites. If we take the sites 1, 2, . . . , $L - 1$ and

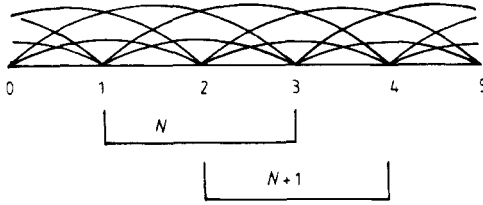


Figure 1. A linear chain with bonds connecting 3rd-nearest neighbours ($L = 3$). Sites (1, 2, 3) and (2, 3, 4) are taken to be the N th and $(N + 1)$ th columns respectively.

L as the N th column, then the $(N + 1)$ th column contains sites 2, 3, . . . , L and $L + 1$. Since each site can be either connected or disconnected to the first column, there are 2^L configurations in each column. The transfer matrix transfers the probability distribution of various configurations in the N th column to the $(N + 1)$ th column. The correlation length ξ is related to the largest non-trivial eigenvalue λ_m of the transfer matrix by the relation (Derrida and Vannimenus 1980)

$$\xi = -1/\ln \lambda_m. \tag{1}$$

From (1), the critical point and the correlation length exponent ν can be found (Zhang and Shen 1982).

To each i th site of the N th column we assign a value m_i which has the value 1 or 0 depending on whether the i th site is connected or disconnected to the first column. So, any configuration in the N th column can be represented by (m_1, m_2, \dots, m_L) . If $P_{m_1 m_2 \dots m_L}(N)$ is the probability of being in the configuration (m_1, m_2, \dots, m_L) , then the transfer matrix is defined by

$$P_{n_1 n_2 \dots n_L}(N + 1) = \sum_{m_1, \dots, m_L = 0, 1} \langle n_1, \dots, n_L | T | m_1, \dots, m_L \rangle P_{m_1 m_2 \dots m_L}(N), \tag{2}$$

where similar notation (n_1, \dots, n_L) is used in the $(N + 1)$ th column. For the site percolation, we have

$$\begin{aligned} \langle n_1, \dots, n_L | T | m_1, \dots, m_L \rangle &= \delta_{n_1 m_2} \dots \delta_{n_{L-1} m_L} (q \delta_{n_L 0} + p \delta_{n_L 1}) && \text{if } (m_1, \dots, m_L) \neq (0, \dots, 0), \\ &= \delta_{n_1 m_2} \dots \delta_{n_{L-1} m_L} \delta_{n_L 0} && \text{if } (m_1, \dots, m_L) = (0, \dots, 0), \end{aligned} \tag{3}$$

where p is the site occupation probability and $q = 1 - p$, and δ is the Kronecker delta.

For the bond percolation, we have

$$\begin{aligned} \langle n_1, \dots, n_L | T | m_1, \dots, m_L \rangle &= \delta_{n_1 m_2} \dots \delta_{n_{L-1} m_L} [q^m \dots q^m \delta_{n_L 0} \\ &\quad + (1 - q^m \dots q^m) \delta_{n_L 1}] && \text{if } (m_1, \dots, m_L) \neq (0, \dots, 0), \\ &= \delta_{n_1 m_2} \dots \delta_{n_{L-1} m_L} \delta_{n_L 0} && \text{if } (m_1, \dots, m_L) = (0, \dots, 0), \end{aligned} \tag{4}$$

where $q_i = 1 - p_i$ and p_i is the occupation probability of the i th-nearest neighbour bond.

Following Domb (1949), we label each configuration (m_1, \dots, m_L) and (n_1, \dots, n_L) of the N th and $(N + 1)$ th columns by numbers m and n respectively, where m and n

Since $A^{(1)}(1, q) = -q$, we find $A^{(L)}(1, q) = -q^L$ and $q_c = 0$ ($p_c = 1$). The value of x of (11) can also be evaluated by a similar procedure and is found to be -1 for any finite L . From (10), we find, for small q , that $\xi = q^{-L}$ and $\nu = L$ (see Klein *et al* 1978). These are the exact results for any finite L .

In the case of bond percolation, $T_{nm}^{(L)}$ is also a duo-diagonal matrix of order 2^L . It is rather difficult to evaluate the determinant $A^{(L)}(1, q_1, \dots, q_L)$ in the general case. However, when all the L th-nearest neighbour bonds have the same occupation probability, the exact critical behaviour can also be obtained by using a symmetric representation as will be shown in the next section.

3. Symmetric representation

Both for the site percolation and the bond percolation with all L th-nearest neighbour bonds having equal occupation probability, a symmetric representation can be used to find the exact critical behaviour. In this representation, we assume that all the configurations having the same number of sites r connecting to the first column have the same probability. We define $P_{\{m\},r}$ in the N th column, as

$$P_{\{m\},r}(N) \equiv P_{m_1 \dots m_L}(N) \quad \text{for all } \sum_{i=1}^L m_i = r \tag{13}$$

where r has the values $0, 1, \dots, L$. If we denote P_r as the total probability of all the configurations having r sites connecting to the first column, we have

$$P_r(N) \equiv \sum_{\substack{m_1, \dots, m_L = 0, 1 \\ m_1 + \dots + m_L = r}} P_{\{m\},r} = \binom{L}{r} P_{\{m\},r} \tag{14}$$

Using the same definition of $P_s(N+1)$ for the $(N+1)$ th column, from (2), we find

$$\begin{aligned} P_s(N+1) &= \sum_{r=0}^L \sum_{\substack{n_1, \dots, n_L = 0, 1 \\ n_1 + \dots + n_L = s \\ m_1, \dots, m_L = 0, 1 \\ m_1 + \dots + m_L = r}} \langle n_1, \dots, n_L | T | m_1, \dots, m_L \rangle \binom{L}{r} P_r(N) \\ &\equiv \sum_{r=0}^L M_{sr}^{(L)} P_r(N) \end{aligned} \tag{15}$$

where $M_{sr}^{(L)}$ is thus the reduced transfer matrix in the symmetric representation having dimensionality $(L+1) \times (L+1)$.

For the case of site percolation, substituting (3) into (15), after some manipulations, we find

$$\begin{aligned} M_{sr}^{(L)} &= \delta_{sr} \quad \text{if } r = 0, \\ &= \delta_{s+1,r} r q / L + \delta_{sr} [r p / L + (L-r) q / L] \\ &\quad + \delta_{s-1,r} (L-r) p / L \quad \text{if } r = 1, 2, \dots, L. \end{aligned} \tag{16}$$

Again, we define $A_{\text{sym}}^{(L)}(\lambda, q)$ as

$$\det(M_{sr}^{(L)} - \lambda \delta_{sr}) \equiv (1 - \lambda) A_{\text{sym}}^{(L)}(\lambda, q). \tag{17}$$

From (16) and (17), $A_{\text{sym}}^{(L)}(1, q)$ can be evaluated exactly and is found to be

$$A_{\text{sym}}^{(L)}(1, q) = F(L)q^L \quad (18)$$

with

$$F(L) = (-1)^L (L-1)! / L^{(L-1)}. \quad (19)$$

It can also be shown that the value of x of (11) is exactly equal to $F(L)$ of (19). Using (10) and (11), we obtain the same results as in the last section; namely, for small q , $\xi = q^{-L}$ and $\nu = L$.

In the case of bond percolation, if we assume that all the L th-nearest neighbour bonds have equal occupation probability, from (4) and (15), with some manipulations, we find

$$\begin{aligned} M_{sr}^{(L)} &= \delta_{sr} && \text{if } r = 0, \\ &= \delta_{s+1,r} \frac{rq^r}{L} + \delta_{sr} \left(\frac{r}{L} (1-q^r) + \frac{(L-r)}{L} q^r \right) \\ &\quad + \delta_{s-1,r} \frac{(L-r)}{L} (1-q^r) && \text{if } r = 1, 2, \dots, L. \end{aligned} \quad (20)$$

Substituting (20) into (17), $A_{\text{sym}}^{(L)}(1, q)$ can again be evaluated exactly and is found to have the form

$$A_{\text{sym}}^{(L)}(1, q) = F(L)q^{L(L+1)/2}, \quad (21)$$

where $F(L)$ is given by (19). In this case the value of x of (11) is again found to be $F(L)$ of (19). From (10) and (11), we obtain, for small q , that $\xi = q^{-L(L+1)/2}$ and $\nu = L(L+1)/2$.

In summary, we have used the transfer-matrix method to find the exact expressions for the correlation length in the critical region for both one-dimensional site and bond percolation systems with bonds connecting L th-nearest neighbours. In the case of site percolation, the correlation length exponent ν is found to be L , consistent with the results obtained by using the generating function method. In the case of bond percolation with all the L th-nearest neighbour bonds having equal occupation probability, ν is found to be $L(L+1)/2$. This confirms the previous prediction by Zhang and Shen (1982).

Recently, we have used the infinitely large cell to cell renormalisation group method proposed by Reynolds *et al* (1980) to treat the bond percolation case (Li *et al*). It is also found that $\nu = L(L+1)/2$, consistent with the results obtained here by using the transfer-matrix method.

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Appendix

For any transfer matrix $T_{nm}(q)$ of order $(J+1)$, the largest non-trivial eigenvalue is given by the largest root of the equation $A(\lambda, q) = 0$, where $A(\lambda, q)$, a determinant

of order J , is defined by

$$\det(T_{nm} - \lambda \delta_{nm}) \equiv (1 - \lambda)A(\lambda, q). \tag{A1}$$

In this appendix, we will show that if $A(1, q)$ has the form bq^K , where b is a constant and K is a positive integer, then the correlation length in the critical region (small q) has the form

$$\xi(q) = xb^{-1}q^{-K}, \tag{A2}$$

where

$$x = \left. \frac{\partial A(\lambda, q)}{\partial \lambda} \right|_{\lambda=1, q=0}. \tag{A3}$$

Since $A(\lambda, q)$ is a polynomial of λ of order J , we can write

$$A(\lambda, q) = \sum_{n=0}^J a_n(q)\lambda^n. \tag{A4}$$

If $\lambda_m(q)$ is the largest root of the equation $A(\lambda, q) = 0$, then $\lambda_m(q)$ must satisfy $\lambda_m(0) = 1$ and

$$A(\lambda_m(q), q) = \sum_{n=0}^J a_n(q)\lambda_m^n(q) \equiv 0. \tag{A5}$$

Expanding $\lambda_m(q)$ in the vicinity of the critical point $q = 0$, we have

$$\lambda_m(q) = 1 + \lambda'_m(0)q + \frac{1}{2}\lambda''_m(0)q^2 + \dots + \frac{1}{n!}\lambda^{(n)}_m(0)q^n + \dots \tag{A6}$$

Taking the total derivatives on both sides of (A5), we have

$$d^n A/d^n q \equiv 0 \quad \text{for all integer } n. \tag{A7}$$

When $n = 1$, we find

$$\sum_{n=0}^J [a'_n(q)\lambda_m^n(q) + na_n(q)\lambda_m^{n-1}(q)\lambda'_m(q)] \equiv 0. \tag{A8}$$

Putting $q = 0$, $\lambda_m(0) = 1$, and using the relation $\sum_{n=0}^J a_n(q) = A(1, q) = bq^K$, (A8) gives

$$\begin{aligned} \lambda'_m(0) &= -x^{-1}bK && \text{if } K = 1, \\ &= 0 && \text{if } K > 1, \end{aligned} \tag{A9}$$

where

$$x = \sum_{n=1}^J na_n(0) \equiv \left. \frac{\partial A(\lambda, q)}{\partial \lambda} \right|_{\lambda=1, q=0}. \tag{A10}$$

If $K > 1$, we can take higher derivatives on both sides of (A8) and put $q = 0$, $\lambda_m(0) = 1$ afterwards. In general, we find the following results:

$$\begin{aligned} \lambda^{(n)}_m(0) &= 0 && \text{if } n < K, \\ &= -x^{-1}bK! && \text{if } n = K. \end{aligned} \tag{A11}$$

From (A6) and (A11), we obtain

$$\lambda_m(q) = 1 - x^{-1}bq^K + O(q^{K+1}). \quad (\text{A12})$$

Substituting (A12) into (1), we finally arrive at (A2) and (A3).

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